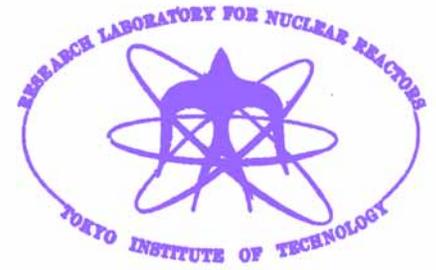
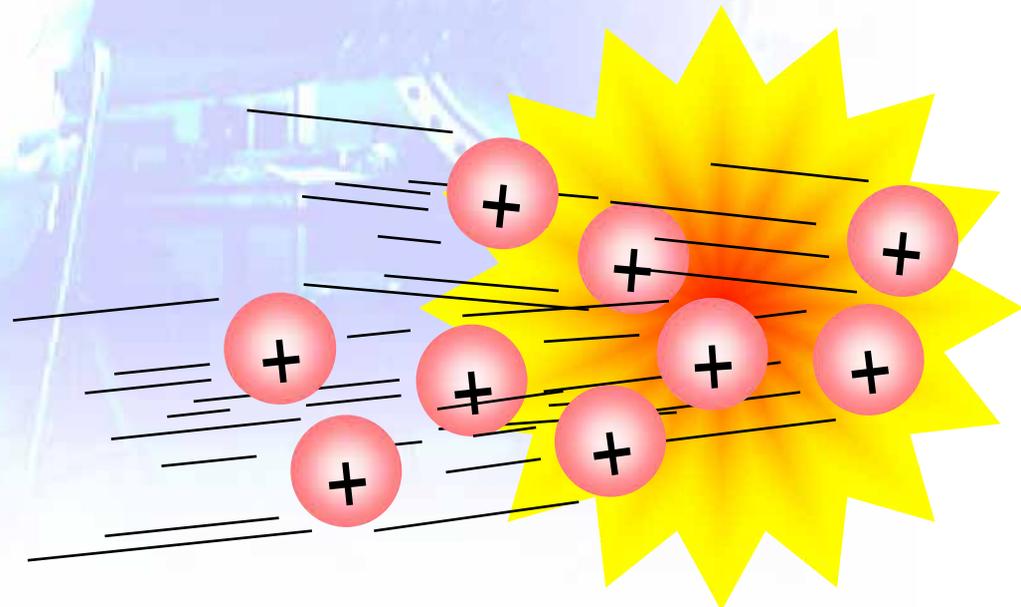


2004 US-Japan Workshop for Heavy Ion Inertial Fusion
10 - 12 June 2004
Princeton Plasma Physics Laboratory



“Beam Stopping Issues and High Energy Density Physics”

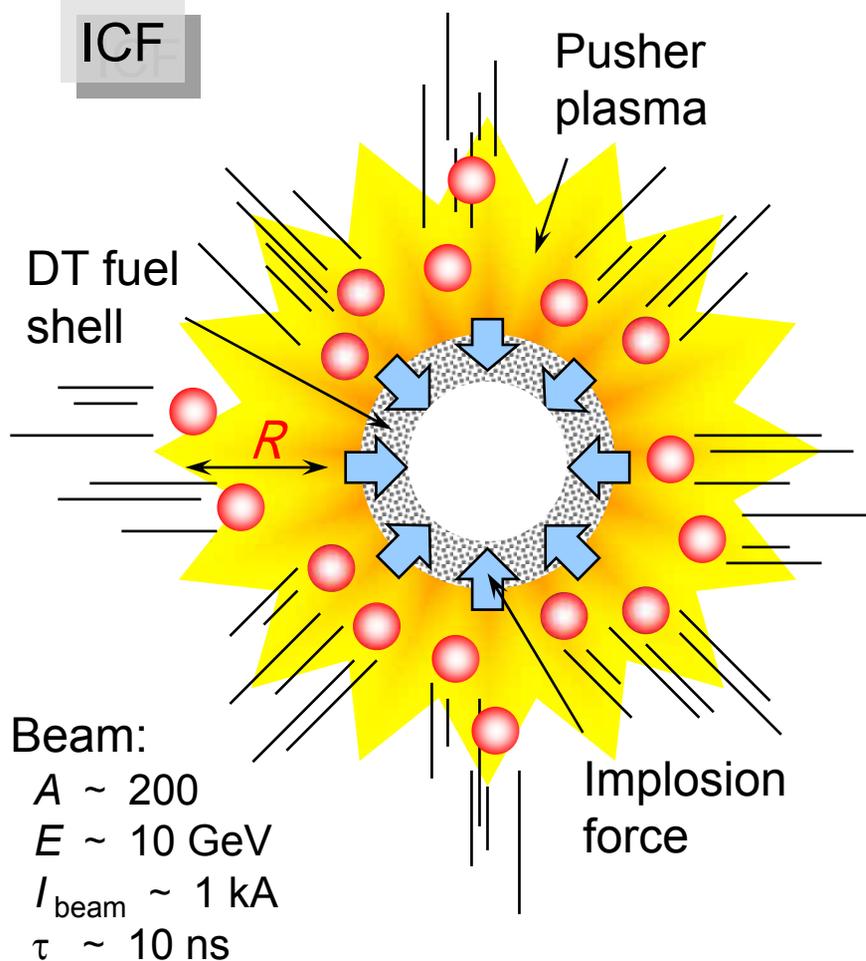
Yoshiyuki Oguri
*Research Laboratory for Nuclear Reactors,
Tokyo Institute of Technology,*



Knowledge on the beam-stopping is of importance to ICF as well as astrophysics.

- e.g., range of energetic ions R in plasmas:

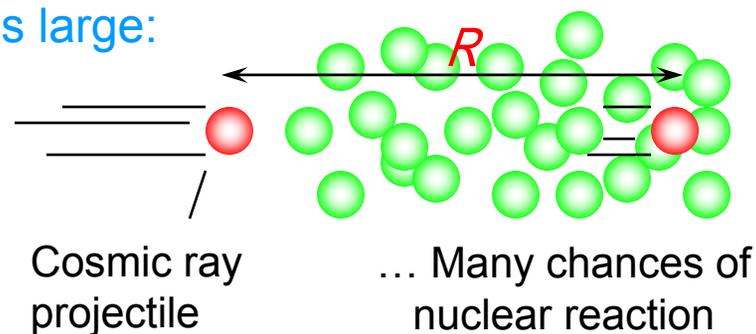
ICF



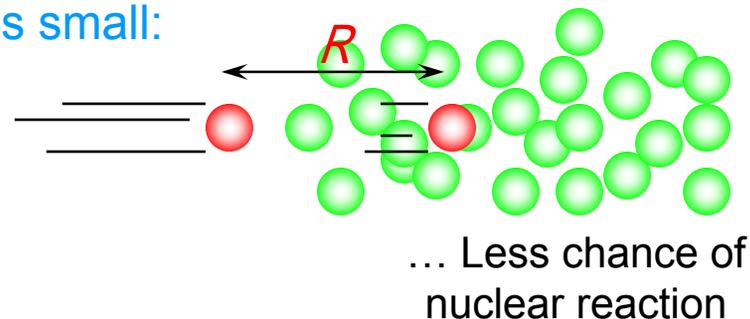
Astrophysics

Example:
Nucleosynthesis in the early solar system

If R is large:

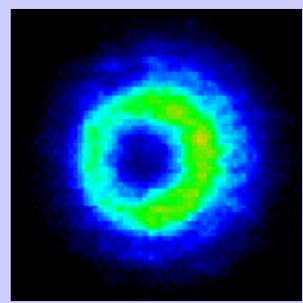
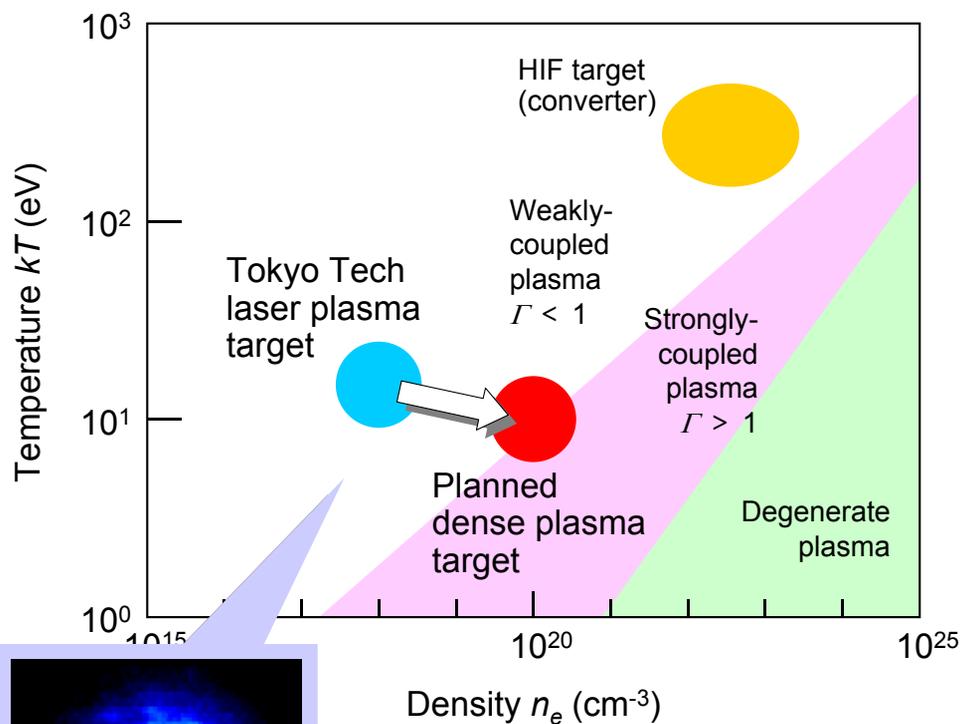


If R is small:



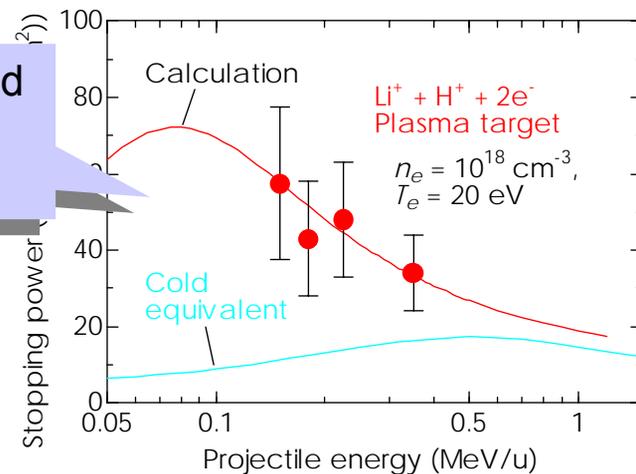
Beam-plasma interaction experiments with dense plasma targets are being planned at RLNR/Tokyo-Tech.

Experiments performed so far using Tokyo-Tech 1.7 MV tandem accelerator:

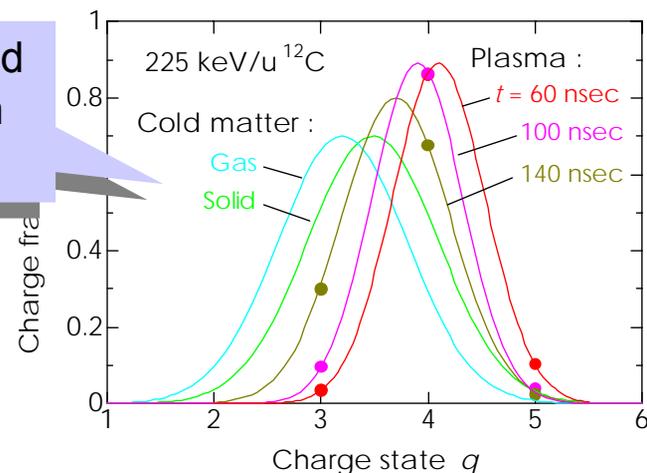


Plasma target $\approx \text{Li}^+ + \text{H}^+ + 2\text{e}^-$,
 $n_e \approx 10^{18} \text{ cm}^{-3}$, $kT \approx 10 \text{ eV}$,

Enhanced $-dE/dx$ in plasmas



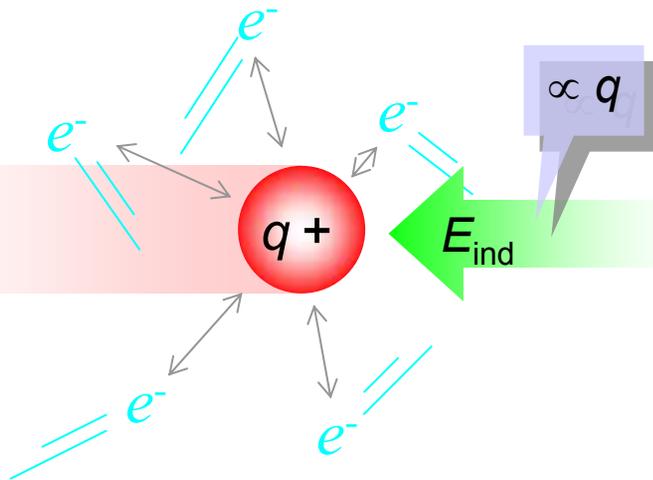
Enhanced charge in plasmas



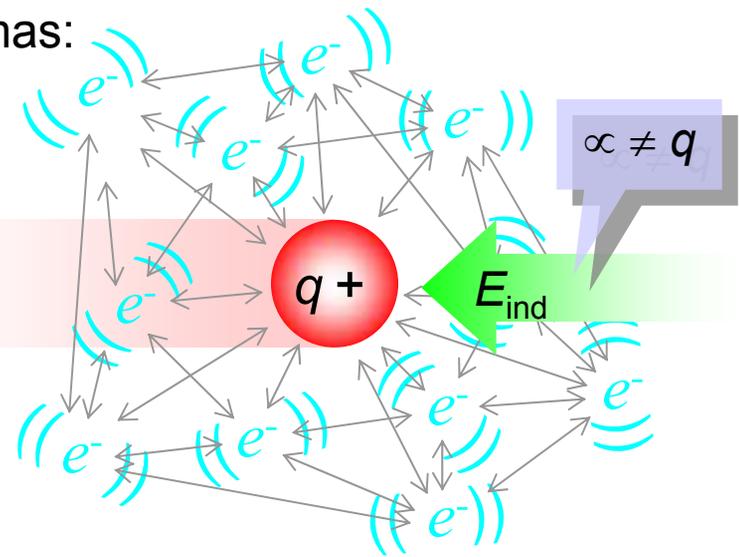
Nonlinear effects are expected for projectile stopping in HIF targets with solid density ($n_e \approx 10^{22} \text{ cm}^{-3}$).

- Dilute hot plasmas ····· Linear stopping:
 - Induced decelerating field $E_{\text{ind}} \propto q$
 - $-dE/dx = qe \times E_{\text{ind}} \quad q \times q = q^2$ (q : projectile charge state)
- Dense cold plasmas ····· **Nonlinear** stopping:
 - Induced decelerating field $E_{\text{ind}} \propto q^m$ ($m < 1$)
 - $-dE/dx = qe \times E_{\text{ind}} \quad q \times q^m = q^{1+m} \equiv q^n$ ($1 < n < 2$)

Dilute plasmas:



Dense plasmas:



Non-ideality of the plasma target is measured by “plasma coupling constant” Γ .

■ Potential energy ϕ and kinetic energy K of plasma particles:

- Mean interparticle distance $\langle r \rangle$, mean Coulomb potential ϕ :

$$\langle r \rangle \equiv \sqrt[3]{\frac{3}{4\pi n_e}}, \quad \phi \equiv \frac{e^2}{4\pi\epsilon_0 \langle r \rangle}$$

- Averaged kinetic energy:

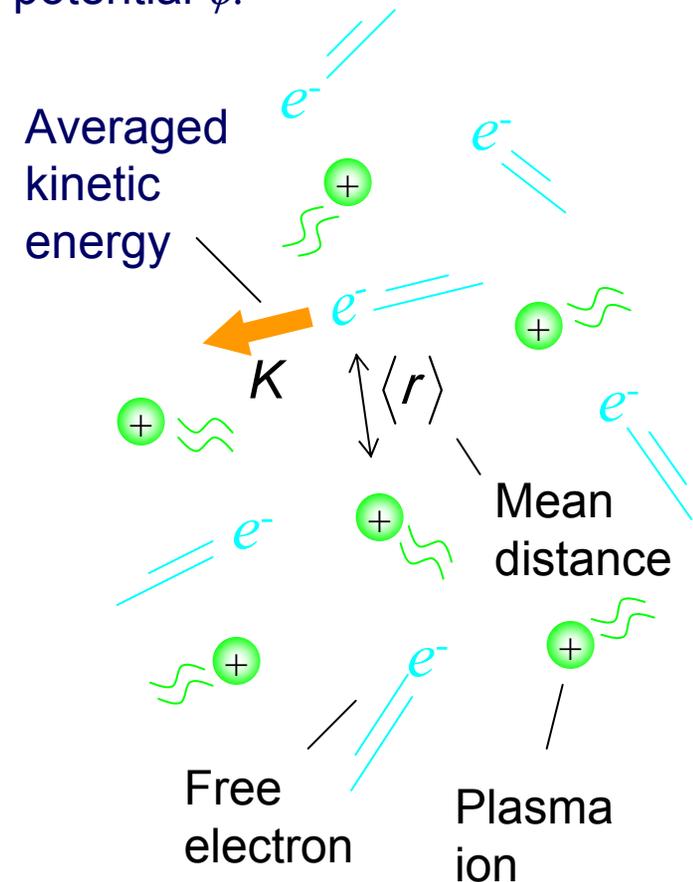
$$K \equiv kT$$

■ Definition of “plasma coupling constant” Γ :

- Ratio of potential energy to kinetic energy:

$$\Gamma \equiv \frac{\phi}{K} = \frac{e^2 \sqrt[3]{4\pi n_e / 3}}{4\pi\epsilon_0 kT}$$

- $\Gamma \approx 0$ “Ideal plasma”
- $\Gamma \neq 0$ “Non-ideal plasma”
- $\Gamma > 1$ “Strongly-coupled plasma”



Strongly-coupled plasmas with $\Gamma \gg 1$ are not always necessary for non-linear stopping experiments.

■ Relationship between the projectile charge q and total charge of electrons which can interact with the projectile:

— Range of Coulomb force by projectile \sim Debye length λ_D : $\lambda_D = \sqrt{\frac{\epsilon_0 kT}{e^2 n_e}}$

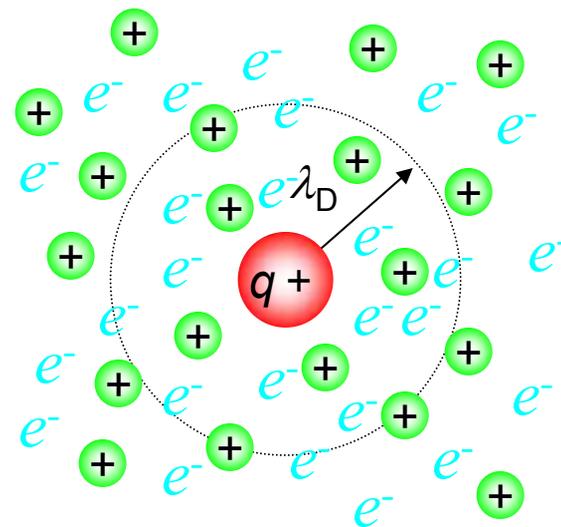
— Number of electrons in a Debye sphere N_D : $N_D = \frac{4\pi}{3} \lambda_D^3 n_e$

— Particle distribution in the plasma is influenced by the projectile charge, if $N_D \ll q$, or,

$$\mathcal{E} \equiv \frac{q}{N_D} = 3^{3/2} q \Gamma^{-3/2} \gg 1$$

■ $\Gamma < 1$ works, if q is very high.

■ Effect of the projectile motion is not included in \mathcal{E} above.



A projectile-plasma coupling constant γ can be defined for projectiles moving in the plasma.

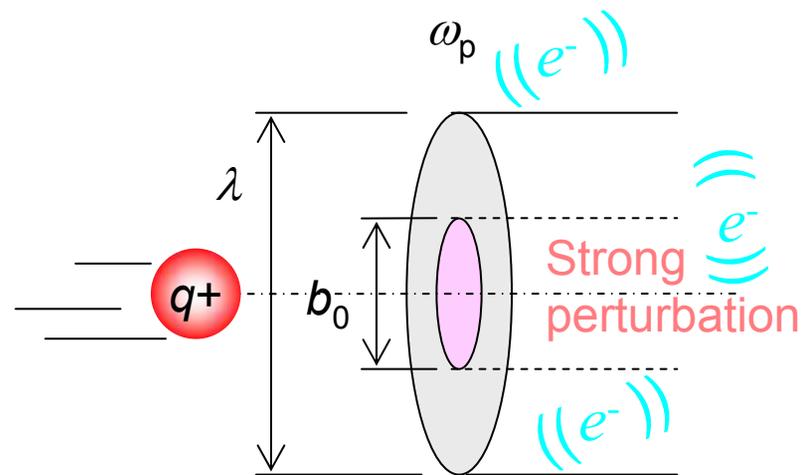
- Perturbations to the plasma electrons are possible only for the collision parameters b smaller than screening length λ :

$$b < \lambda = \frac{\langle v_r \rangle}{\omega_p}, \quad \langle v_r \rangle = v_{th} \sqrt{1 + \left(\frac{v_{proj}}{v_{th}} \right)^2}$$

— $\langle v_r \rangle$: averaged relative energy

- If b is smaller than the classical collision diameter b_0 , the perturbation is strong enough to induce nonlinear effects:

$$b < b_0 \equiv \frac{qe^2}{4\pi\epsilon_0 m \langle v_r \rangle^2}$$

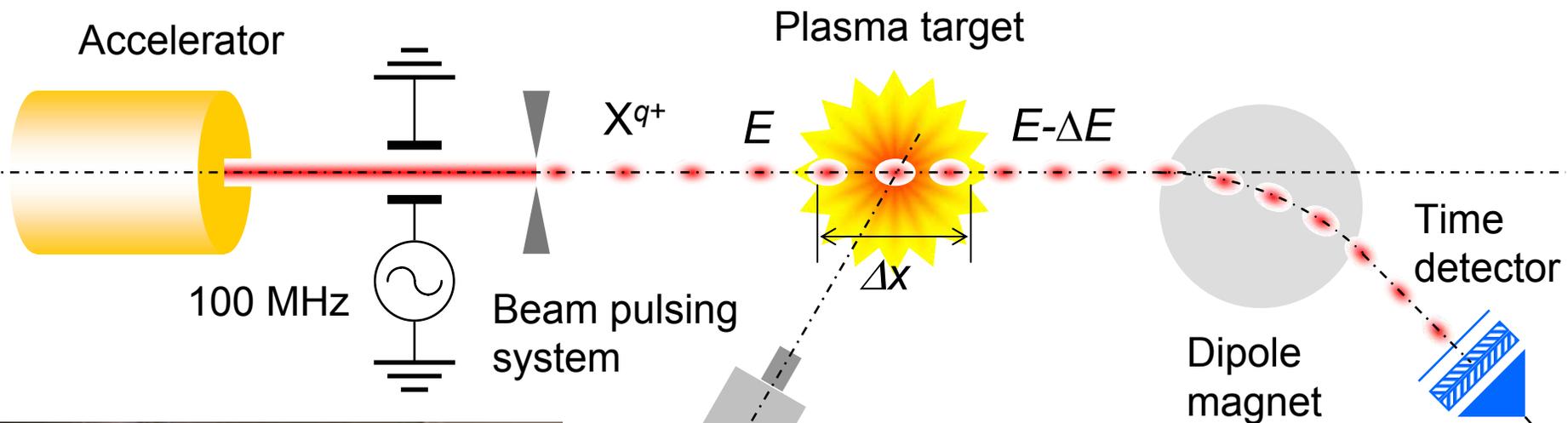


- The projectile-plasma coupling strength is estimated by the critical ratio $\gamma \equiv b_0/\lambda$:

$$\gamma \equiv \frac{b_0}{\lambda} = \frac{qe^2 \omega_p}{4\pi\epsilon_0 m \langle v_r \rangle^3} = \frac{\sqrt{3}q\Gamma^{3/2}}{\left\{ 1 + \left(\frac{v_{proj}}{v_{th}} \right)^2 \right\}^{3/2}}$$

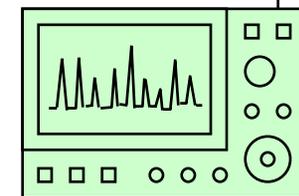
We are looking for appropriate experimental parameters to observe nonlinear effects.

■ Energy loss measurement by Time-Of-Flight method:



Plasma diagnostics

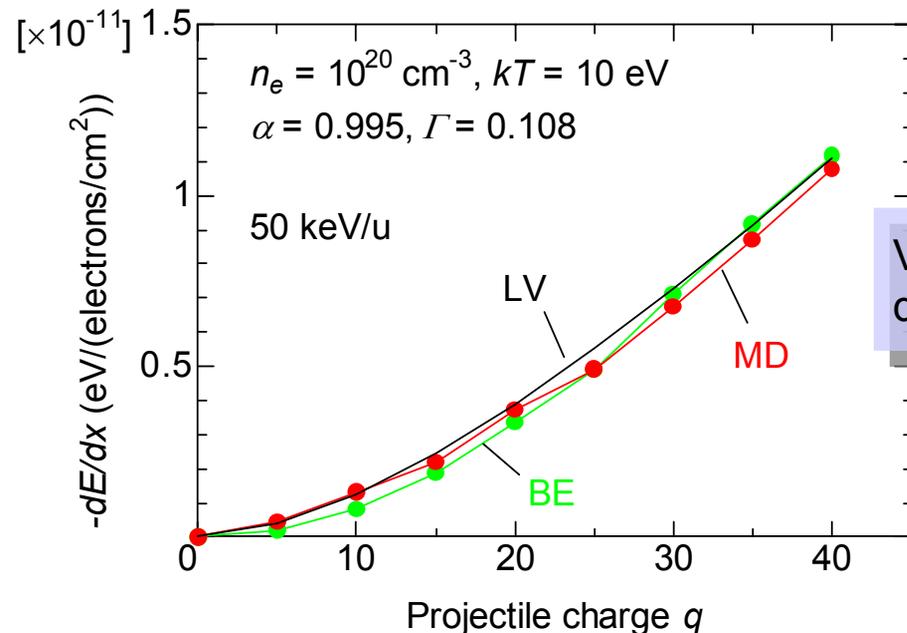
Target : $n_e, kT, \Delta x$?
 Projectile: E, q ?



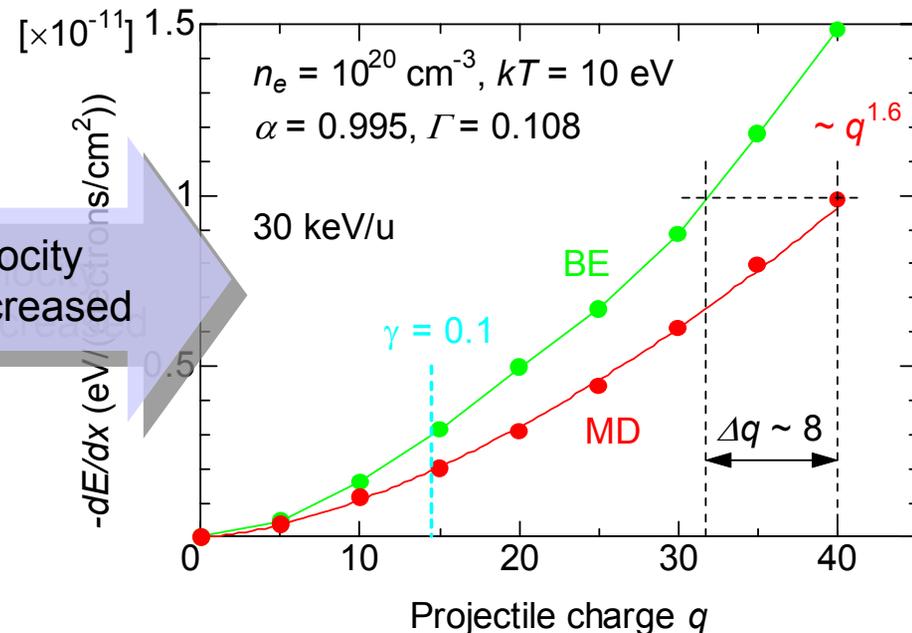
Storage oscilloscope

Nonlinear effect is remarkably increased by slightly decreasing the projectile velocity.

- 30 keV/u is acceptable, although lower projectile energies are not preferable as practical experimental conditions.
- $q > \approx 15+$ may be necessary to clearly observe the nonlinear effects.
- For $q = 40+$, the decrease of the projectile effective charge is ≈ 8 .
 - At least 8 electrons are responsible for the screening ?

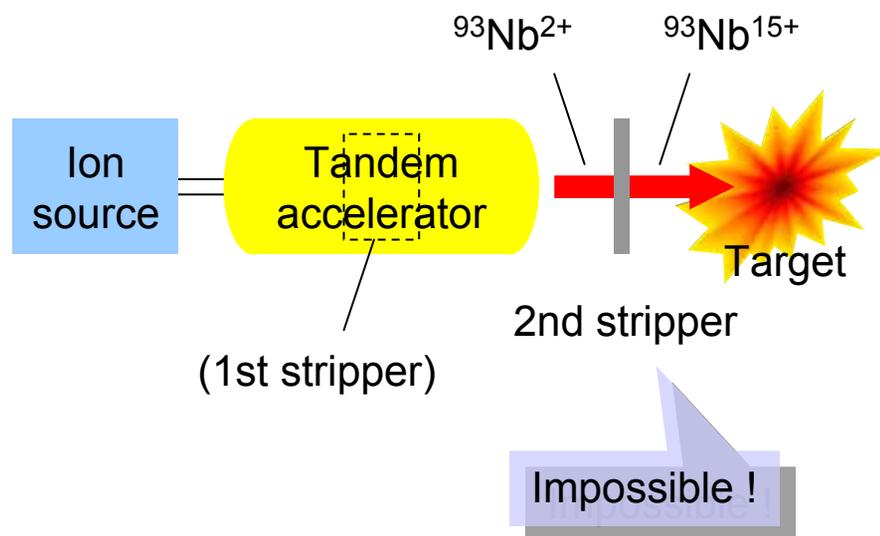
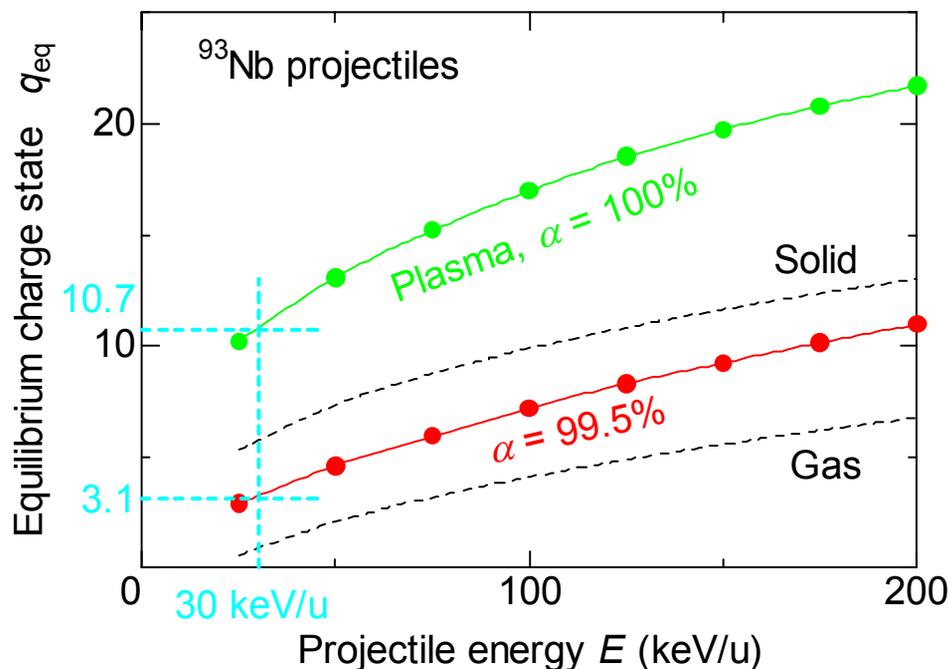


Velocity decreased



Experiments using a 30 keV/u ^{93}Nb beam from the tandem accelerator is very difficult.

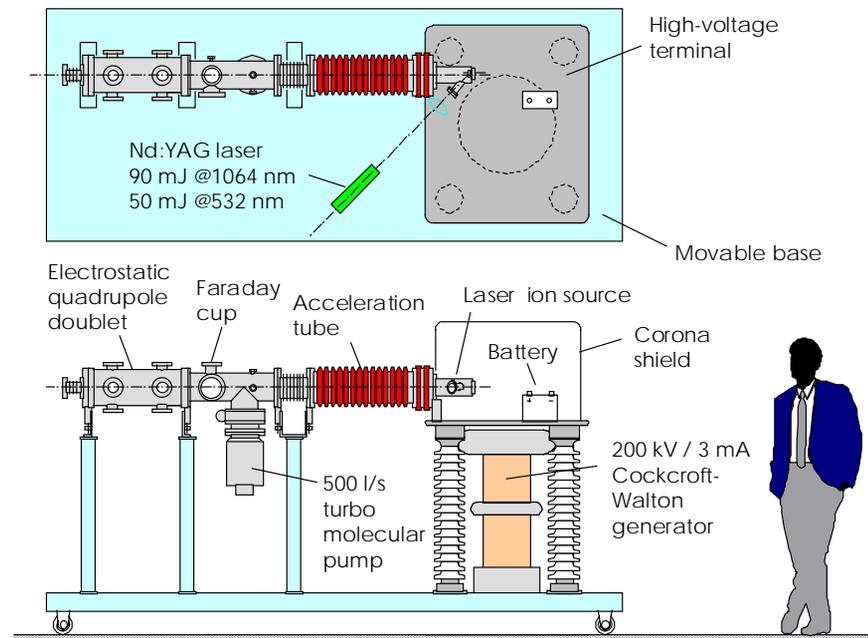
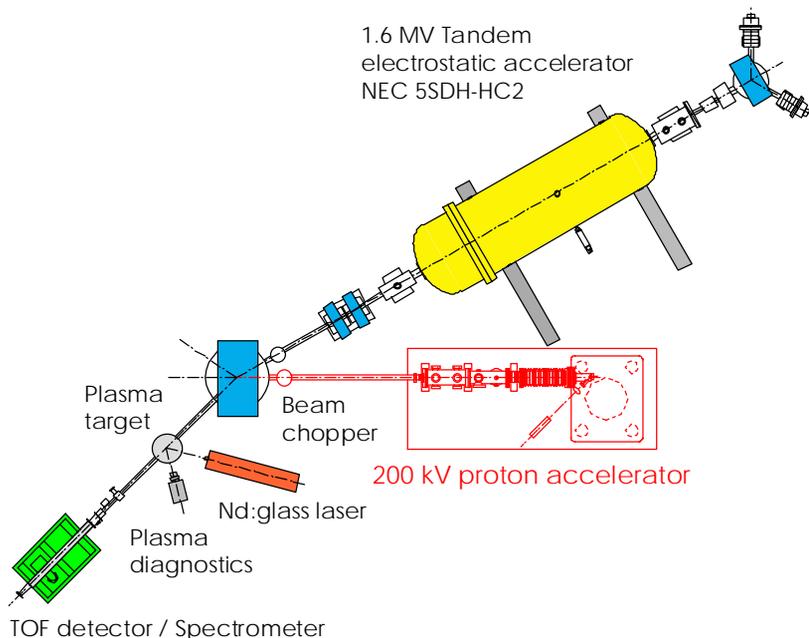
- Charge state of the 30 keV/u ^{93}Nb beam from the accelerator is only $q = 2+$.
→ second stripper is necessary.
- At 30 keV/u, $q \approx 15+$ is not available by ordinary stripping processes:
 - e.g., stripping of $_{41}\text{Nb}$ by C-foil $q \approx 7$
 - Also plasma stripper does not work very well.



As an alternative, a single-ended low-energy machine with a source of highly-charged ions is considered.

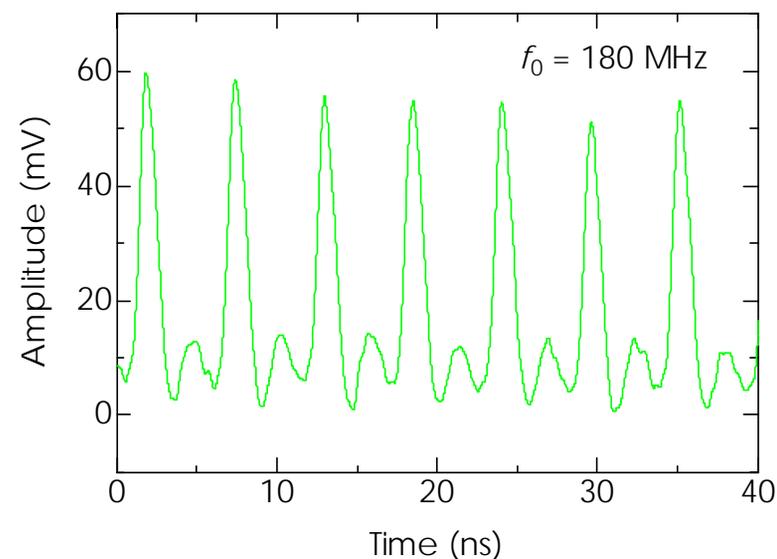
- Originally designed for experiments using 100-200 keV protons for the measurement of Coulomb logarithm $\ln \Lambda$:
 - Laser Ion source (YAG- 2ω , 50 mJ, 5 ns)
 - 180 MHz reentrant resonator cavity for beam bunching
- The beam is injected into the existing beam line.

$$-dE/dx \propto q_{\text{eff}}^2 \ln \Lambda$$



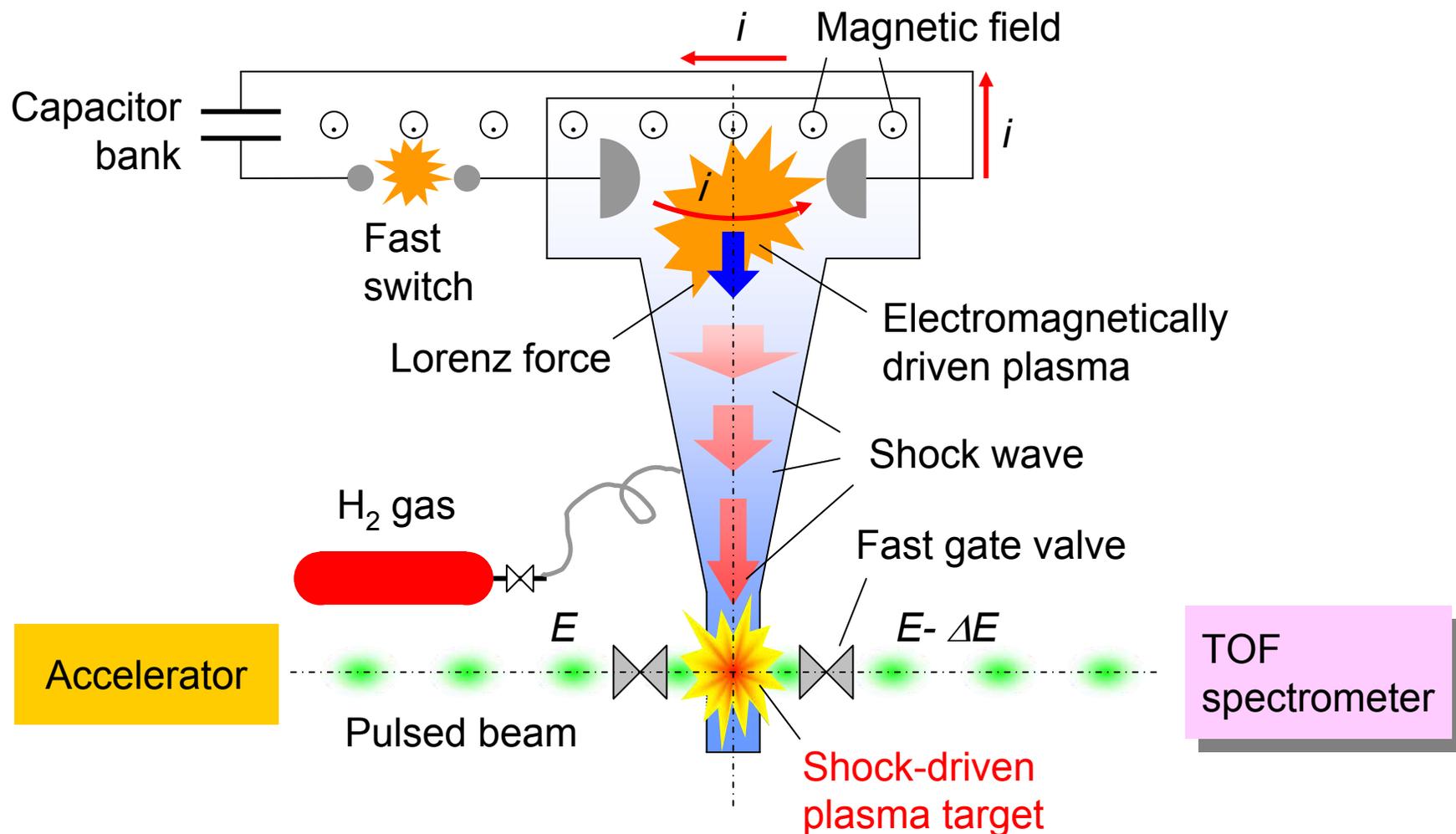
If the small YAG laser is changed to an existing glass laser, highly-charged projectiles will be obtained.

- By using the existing Nd:glass laser ($1.05 \mu\text{m}$, 4 J, 30 ns), irradiation of $\approx 10^{12} \text{ W/cm}^2$ is possible.
- If heavy ions with $q = 15+$ are available, experiments with the projectile energies up to $200 \text{ kV} \times 15e = 3.0 \text{ MeV}$ are possible.
 - Cf. at CERN, 70 mA $^{181}\text{Ta}^{20+}$ by a 10^{12} W/cm^2 CO_2 laser.
- The machine is under the test of beam transport to the plasma target.



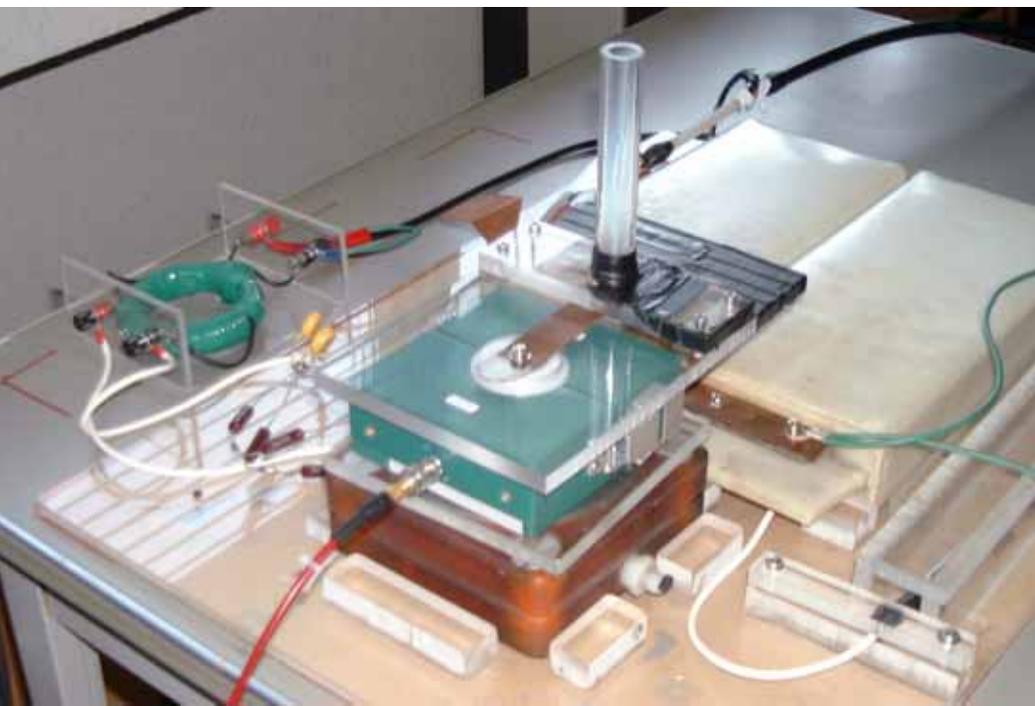
An electromagnetically-driven shock tube is being developed to produce cold dense plasma targets.

- Projectile energy loss is measured by a TOF method:



A prototype of the electromagnetic shock tube is being tested with air at the atmospheric pressure.

- Parameters of preliminary experiments:
 - Tube diameter / length = 14 mm / 140 mm
 - Capacitor bank $C = 1 \mu\text{F}$, $V = 20 \text{ kV}$
 - Energy = 200 J
 - Time constant of discharge = $2.5 \mu\text{s}$



$V = 10 \text{ kV}$,
Exposure = $1 \mu\text{s}$

